

**ON THE REDUCTION OF CERTAIN INTEGRAL EQUATIONS
OF THE FIRST KIND IN THE THEORY OF ELASTICITY
AND HYDRODYNAMICS TO EQUATIONS OF THE SECOND KIND**

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It is shown here that the solution of certain integral equations of the first kind which are encountered in applications reduce to solution of two equations: 1) the simplest equation of the first kind, the Abel integral equation; and 2) the equation of the second kind.

1. We shall examine the equation of the theory of a slender wing ([1], p. 80)

$$\int_{\alpha_0}^{\alpha_1} u(\alpha) \left(\frac{2a + \alpha_1 - \alpha}{\alpha_1 - \alpha} \right)^{1/2} d\alpha = \int_{-a}^a (v_{n0} - v_n) \left(\frac{a - \xi}{a + \xi} \right)^{1/2} d\xi \quad \left(\begin{array}{l} \alpha_1 > \alpha_0 \\ a > 0 \end{array} \right) \quad (1.1)$$

with the unknown function $u(\alpha)$. For $a = \text{const}$, the kernel of this equation depends only on the difference of the arguments, and Eq. (1.1) can be solved in closed form with the aid of Mellin's formula [2]; however, as noted in [1], this approach is not effective because of a computational difficulty.

Replacing the right-hand side of (1.1) by an arbitrary function, we rewrite (1.1) in the form

$$\int_{\alpha_0}^x u(\tau) \left(\frac{2a + x - \tau}{x - \tau} \right)^{1/2} d\tau = f(x, a) \quad (x > \alpha_0) \quad (1.2)$$

or

$$\sqrt{2a} \int_{\alpha_0}^x \frac{u(\tau) d\tau}{\sqrt{x - \tau}} = \int_{\alpha_0}^x \frac{\sqrt{2a} - \sqrt{2a + x - \tau}}{\sqrt{x - \tau}} u(\tau) d\tau + f(x, a) \quad (1.3)$$

Inverting (1.3) as an Abel equation we have

$$u(x) + \frac{1}{\pi \sqrt{2a}} \frac{d}{dx} \int_{\alpha_0}^x \frac{d\tau}{\sqrt{x - \tau}} \int_{\alpha_0}^{\tau} \frac{\sqrt{2a + \tau - \xi} - \sqrt{2a}}{\sqrt{\tau - \xi}} u(\xi) d\xi = f_1(x, a) \quad (1.4)$$

$$f_1(x, a) = \frac{1}{\pi \sqrt{2a}} \frac{d}{dx} \int_{\alpha_0}^x \frac{f(\tau, a)}{\sqrt{x - \tau}} d\tau$$

Changing the order of integration in (1.4) and differentiating under the integral sign, we obtain

$$u(x) + \frac{1}{4a\pi} \int_{\alpha_0}^x M \left(\frac{x - \xi}{2a} \right) u(\xi) d\xi = f_1(x, a) \quad (1.5)$$

$$M \left(\frac{x - \xi}{2a} \right) = 4a \frac{d}{dx} \int_{\xi}^x \left[\left(1 - \frac{\tau - \xi}{2a} \right)^{1/2} - 1 \right] \frac{d\tau}{\sqrt{x - \tau} \sqrt{\tau - \xi}} =$$

$$= \int_0^1 \left(\frac{s}{1-s} \right)^{1/2} \frac{ds}{\sqrt{1 + 1/2s(x-\xi)/a}} = \frac{\pi}{2} F \left(\frac{1}{2}, \frac{3}{2}; 2; -\frac{x-\xi}{2a} \right) \tag{1.6}$$

The Volterra Eq. (1.5) has a kernel of difference type and can, therefore, be solved in closed form like the original equation. However, unlike (1.2), the solution of Eq. (1.5) can be obtained simply by the method of successive approximations.

We note that the indicated method of reduction of Eq. (1.2) to Eq. (1.5) is also applicable in the case $a \neq \text{const}$, but then the kernel of Eq. (1.5) will not depend only on the difference $(x-\xi)$.

By means of the change of variables in (1.6)

$$s = 1 - \frac{y^2}{k^2}, \quad k = \left(\frac{x}{1+z} \right)^{1/2} = \left(\frac{x-\xi}{2a+x-\xi} \right)^{1/2} \quad (0 < k < 1)$$

the kernel $M(z)$, $z = \frac{1}{2}(x-\xi)/a$ can be reduced to the form

$$M(z) = \frac{2}{\sqrt{z}} \int_0^k \left[\frac{1}{1-y^2} \left(1 - \frac{y^2}{k^2} \right) \right]^{1/2} dy = \frac{2}{\sqrt{z}} E \left(\sin^{-1} k, \frac{1}{k} \right)$$

where $E(\varphi, 1/k)$ is the elliptic integral of the second kind. We may also express $M(z)$ in terms of complete elliptic integrals of the first and second kinds

$$M(z) = \frac{2\sqrt{1+z}}{z} \left[E(k) - \frac{K(k)}{1+z} \right]$$

(see 8.111, 8.112 or 3.169, 9 in [9]).

2. The method of separation of the Abelian part of an equation which has been illustrated by the example of Eq. (1.2) can be generalized to other equations. Besides the trivial generalization to the case of the kernel

$$\left(\frac{2a+x-\tau}{x-\tau} \right)^\mu \quad (0 < \mu < 1)$$

in (1.2), this method can be successfully applied to Eq.

$$\int_a^b \frac{c_1 + c_2 \text{sign}(x-\tau)}{|x-\tau|^\mu} \varphi(\tau) d\tau + T\varphi = f(x) \tag{2.1}$$

where c_1 and c_2 are functions which depend only on x or only on τ , and the kernel of the integral operator T either has no singularity whatever on the diagonal $x=\tau$, or else has a singularity of power type with exponent smaller than μ .

Various special cases of this equation (which for $T=0$ has been called the generalized Abel equation [3]) have important applications in the theory of elasticity [4], the theory of creep and the theory of plasticity [5 to 7]. Equation (2.1) with $T=0$ was first solved by Sakaliuk [3 and 8] by the method of analytic continuation. This equation was investigated by the author of the present paper using another method. We indicate here only that by separating the Abelian part

$$\int_a^x \frac{c_1 + c_2}{(x-\tau)^\mu} \varphi(\tau) d\tau$$

out of (2.1) and applying to (2.1) the operator $I_{ax}^{-(1-\mu)}$, which is the inverse of the operator

$$I_{ax}^{1-\mu} \varphi \equiv \frac{1}{\Gamma(1-\mu)} \int_x^x \frac{\varphi(\tau) d\tau}{(x-\tau)^\mu}$$

we reduce Eq. (2.1) to a singular equation with a kernel of Cauchy type. The operator $I_{ax}^{-(1-\mu)}$ must be applied to (2.1) on the left if the coefficients C_1 and C_2 depend on T and on the right if they depend on x . The following identities have been used:

$$\begin{aligned} I_{ax}^{-\nu} I_{\tau\beta}^{\nu} \varphi &\equiv \cos(\nu\pi) \varphi(x) + \frac{\sin \nu\pi}{\pi} \frac{1}{(x-a)^\nu} \int_a^\beta \frac{(\tau-a)^\nu \varphi(\tau) d\tau}{\tau-x} \\ I_{x\beta}^{\nu} I_{a\tau}^{-\nu} \Phi &\equiv \cos(\nu\pi) \Phi(x) + \sin(\nu\pi) \frac{(b-x)^\nu}{\pi} \int_a^\beta \frac{\Phi(\tau) d\tau}{(b-\tau)^\nu \tau-x} \\ \left(I_{x\beta}^{\nu} \varphi &\equiv \frac{1}{\Gamma(\nu)} \int_x^\beta \frac{\varphi(\tau) d\tau}{(\tau-x)^{1-\nu}}, \nu = 1-\mu \right) \end{aligned} \quad (2.2)$$

Equations (2.2) allow us to arrive at an equation with Cauchy kernel which for $T=0$ is solvable in closed form. They also make it possible to eliminate (where necessary) the singular integrals with the aid of which the solution of the equation with Cauchy kernel is expressed. In conclusion we note that Eq. (1.2) is a particular case of Eq. (2.1) with $c_1 = c_2 = \sqrt{1/2} a$, $\mu = 1/2$ and

$$T\varphi \equiv \int_a^x \frac{\sqrt{2a+x-\tau} - \sqrt{2a}}{\sqrt{x-\tau}} \varphi(\tau) d\tau \equiv \int_a^x \frac{\sqrt{x-\tau}}{\sqrt{2a} + \sqrt{2a+x-\tau}} \varphi(\tau) d\tau$$

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